Trapping and Escape of Buoyant Plumes in Stratified Water

 Rich McLaughlin Center for Interdisciplinary Applied Mathematics and Joint Fluids Lab, UNC Chapel Hill

Close Collaboration with **Roberto Camassa** and **Brian White**

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Internal Splash

- 1.0396 g/ cc
- **Stratified**
- Sphere heavier than fluid at all depths
- Fast Playback
- Re=300

Abaid,Adalsteinsson,Agyapong,McL Phys of Fluids 2004

Srdic-Mitrovic, Mohamed Fernando, JFM 1999 (no bounce)

1.0385g/CC Srdic-Mitrovic, Mohamed,Fernando, JFM 1999 (no bounce) For keynote with movies: See pink-lady.amath.unc.edu/~rmm/sost

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Left: low speed untreated oil, right: high speed+dispersant

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Slide: from Harvey Seim, **Marine Sciences, UNC**

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JAG, 2010

Integral Models: Morton, Taylor, Turner, 1953

The hight where the oil got trapped is determined by the location where *w*⁰ goes

TRACK SINGULARITY

λ2

 T rac \mathcal{L}_R Singularity is a set of \mathcal{L}_R . The set of \mathcal{L}_R is a set of \mathcal{L}_R

CHUNG-NAN TZOU

$$
\frac{d(b^2w)}{dz} = 2\alpha bw,
$$
\n
$$
\frac{d(b^2w^2)}{dz} = 2g\lambda^2b^2\theta,
$$
\n
$$
\frac{d(b^2w\theta)}{dz} = -\frac{1+\lambda^2}{\lambda^2}\frac{d\epsilon}{dz}b^2w
$$

 ϵ -- ambient stratification *dz* (1) = 2α*bw,* to 0, which is the singular point of the speed function *w*(*z*). and the system of

The parameters used are:

dz (1) = 2α*bw,*

[−]¹ ⁺ ^λ²

d(*b*²*w*)

d(*b*²*w*²)

d(*b*²*w*θ)

d(*b*²*w*)

to 0, which is the singular point of the speed function *^w*(*z*). and the system of

dz (1) ⁼ ²α*bw,*

g (*cm/sec*²) λ α ρ*^b* ρ*^t* ρ*^j b*⁰ *w*⁰ θ⁰ A, C ⁻⁻entrainment, mixing coefficients \mathbf{b} **b**, w, θ $\mathsf{b}, \mathsf{w}, \ \theta$ **d** θ , w, θ λ, α dzie *b*² (3) *w.*
 dzie biologiczne bi -- jet radius, center speed, density --entrainment, mixing coefficients **d**_{*d*} $\frac{1}{2}$ (2) $\frac{1}{2}$ (3) $\frac{1}{2}$ λ ,(*dz ^b*² (3) *w. dz* (1) = 2α*bw,*

^λ²))(^ρ*b*−ρ*^j*

Miscible Limit: Critical Escape Height For Buoyant Jets

$$
L = L_0 \int_1^A \frac{ds}{\sqrt{s^{5/4} + \epsilon - 1}}
$$

$$
\epsilon = \frac{5(1+\lambda^2)(\Delta\bar{\rho})r_0g}{16\sqrt{2}\alpha w_0^2} \quad A = (1 + \epsilon(\frac{\theta_0^2}{\theta_f^2} - 1))^{4/5}
$$

$$
\theta_0 = \frac{(1+\lambda^2)}{\lambda^2} \Delta\bar{\rho} \quad \theta_f = \frac{\rho_b - \rho_t}{\rho_b} \quad \Delta\bar{\rho} = \frac{\rho_b - \rho_j}{\rho_b}
$$

$$
L_0 = \left(\frac{5r_0w_0^2}{16\sqrt{2}g(1+\lambda^2)\alpha(\Delta\bar{\rho})}\right)^{1/2}
$$

Adalsteinnson. Camassa, Falcon, Lin, McLaughlin, Mertens, Nenon, Smith, Walsh, Watson, White, to appear: "Monitoring and Modeling the Deepwater Horizon Oil Spill: A Record-Breaking Enterprise, AGU Monograph Series

Asymptotics-- Camassa, McL, Tzou, Zhao, in prep

Current Events: Ocean Carbon Pump

Plume Destabilization

Figure 3. Time series showing timescale of plume instability. Top: OSW 4:3:2, $t = 30$, 450, 870, 900, 1800, 3600, 7200 sec. Bottom: OSW 4:3:17, *t* = 30, 450, 900, 1800, 3600, 7200, 14400 sec. Notice the onset of instability in the top row, first evident at $t = 870$ sec.

Top: 48 cm travel, Bottom 15cm Left to right: increasing flow rate

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Stratified Vortex Rings

Camassa, McL, **Keith Mertens**, D. Nenon, C. Smith, C. Viotti, in prep

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DNS: Varden (A. Almgren LBL Code) modified by Claudio Viotti

 mesh:256x256x1024, parallel on 256 processors run time: 6 hours periodic x-y, slip wall velocity lids no flux bc for scalar

Experiment

3/5 power law

0.02 0.03 0.04 0.05 0.06 0.07 0.08 0.09 0.10 0.11

Figure 3: Critical length scale for total trapping of miscible buoyant jets and vortex rings: Left

panel, buoyant jet data, (un-fit) scaling prediction, and exact hypergeometric formula, length scaling prediction, and exact hypergeometric formula, length scaling prediction, and exact hypergeometric formula, length scal

Exact

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Conclusions and Future:

Buoyant plume formation in stratification Trapping timescales vary with mixture Plume destabilization may occur Internal waves-- larger scale experiments Inflow full DNS CFD marginally resolvable

UNC Joint Fluids Lab, Chapman Hall Level B, rm B02

Why Did Huge Oil Plumes Form After the Gulf Spill

Weird

Nev

Dispersants broke oil into micro-droplets suspended by equally dense water.

Oil jets pre-mixed with soap are fired into layered fluid, mimicking the spreading of a Gulf oil plume. Photograph by Steve Harenberg, Rich McLaughlin, Johnny Reis, William Schlieper, Will Owens, Brian White, UNC Joint Fluids Laboratory and UNC Center for Interdisciplinary Applied Mathematics Roberto Camassa/The University of North Carolina